

to all the planes and all the directions in them, i.e., within the limits

$$\begin{aligned} 0 \leq \alpha \leq 2\pi \quad -\pi/2 \leq \varphi \leq \pi/2 \\ -\pi/2 \leq \beta \leq \pi/2 \end{aligned} \quad (9)$$

where

$$Q_1 = 1/4\pi \quad Q_2 = 1/4\pi^2 \quad (10)$$

On the basis of this and relations (8), Eqs. (4) assume the form

$$\begin{aligned} \sigma_{ij} = \frac{a}{4\pi} \int_0^{2\pi} d\alpha \int_{-\pi/2}^{\pi/2} d\varphi [\epsilon_{nn} l_{in} l_{jn}] + \frac{b}{4\pi^2} \int_0^{2\pi} d\alpha \times \\ \int_{-\pi/2}^{\pi/2} d\varphi \int_{-\pi/2}^{\pi/2} d\beta [\epsilon_{nm} (l_{in} l_{jm} + l_{im} l_{jn})] \end{aligned} \quad (11)$$

Taking into account Eqs. (6), after integration we get

$$\sigma_{ij} = \lambda \theta \delta_{ij} + 2\mu \epsilon_{ij} \quad (12)$$

that is, the ordinary generalization of Hooke's law, where

$$\lambda = (a + b)/15 \quad \mu = (2a + 3b)/30 \quad (13)$$

On the basis of known equations connecting Lamé's constants with Young's modulus  $E$  and Poisson's ratio  $\nu$  we get

$$E = \frac{a}{3} \left[ \frac{2a + 3b}{4a + b} \right] \quad \nu = \frac{a - b}{4a + b} \quad (14)$$

Whence it follows that

$$a = 3E/(1 - 2\nu) = 3K \quad (15)$$

Here  $K$  is the modulus of bulk compression.

Suppose that the yield stress, measured in a uniaxial tension test, is  $\sigma_H$ , whereas the theoretical stress, calculated from conditions of rupture of the atomic planes, is  $\sigma_T$ . The maximum tearing stresses are experienced by atomic planes perpendicular to the direction in which the tension is applied. By virtue of relations (6)

$$\sigma_{nn} = \sigma_T = a\epsilon_{nn} = a\epsilon \quad (16)$$

where  $\epsilon$  is the axial deformation. On the other hand,  $\epsilon = \sigma_0/E$ ; but then on the basis of Eq. (16) we have

$$\sigma_H = 3(1 - 2\nu)\sigma_T \quad (17)$$

From this it follows that, as  $\nu$  approaches  $\frac{1}{2}$  (incompressible body), the observed rupture stress may be much less than the theoretical stress. Although (in view of the asymmetry of the method with respect to stresses and deformations) this cannot be regarded as a new explanation of the sharp difference between the theoretical and observed rupture stresses, a more critical approach is necessary to existing calculations of theoretical stress with regard to interactions not coinciding with the direction of rupture.

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## Cycles about a Singular Point of Nodal Type

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\*In a differential equation system in the plane

$$dx/dt = P(x, y) \quad dy/dt = Q(x, y)$$

where  $P$  and  $Q$  are polynomials, with  $P(0,0) = Q(0,0) = 0$ , there may exist closed trajectories enclosing the singular point at the origin. If the degrees of  $P$  and  $Q$  are at most two, and the origin is a nodal point, it is shown that there can be no closed trajectory enclosing the origin. Sufficient conditions for a nodal point to be acyclic (i.e., not to be enclosed by a closed trajectory) are developed.

IT is known that for a system of differential equations of the form

$$\frac{dx}{dt} = \sum_{i+j=1}^n a_{ij} x^i y^j \quad (1)$$

$$\frac{dy}{dt} = \sum_{i+j=1}^n b_{ij} x^i y^j$$

( $a_{ij}, b_{ij}$  constant) that when  $n = 1$  (a linear system) a singular point of nodal type cannot be enclosed by a closed trajectory.

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This is possible for  $n \geq 3$ . For example, in the system

$$\begin{aligned} dx/dt &= -y + ax(x^2 + y^2 - 1) \\ dy/dt &= x + by(x^2 + y^2 - 1) \end{aligned}$$

with  $ab \geq -1$ ,  $(a - b)^2 \geq 4$ , the singular point at  $(0,0)$  is a node, and it is enclosed by the closed trajectory  $x^2 + y^2 = 1$ .

In the present note we shall formulate sufficient conditions that a singular point be acyclic,<sup>†</sup> and shall show that in the system (1) with  $n = 2$  a singular point of nodal type cannot be enclosed by a closed solution.

Consider the system of differential equations

$$d\rho/dt = F(\rho, \theta) \quad d\theta/dt = \Phi(\rho, \theta) \quad (2)$$

(where  $\rho$  and  $\theta$  are polar coordinates and  $F(0, \theta) \equiv 0$ ), for which conditions for existence and uniqueness of solutions are fulfilled throughout the entire plane.

<sup>†</sup> By an acyclic singular point we shall mean a singular point not enclosed by a closed trajectory.

**Lemma:** Let the graph of the equation  $\Phi(\rho, \theta) = 0$  have a branch  $\Gamma$ , a portion of which,  $\Gamma_1$ , extends to infinity. Let  $\Gamma_1$  be contained in the sector  $\theta_1 \leq \theta \leq \theta_2$  where  $\theta_2 - \theta_1 < 2\pi$ . Let  $\Phi(\rho, \theta)$  change sign as  $\Gamma_1$  is crossed, and let there be no other branch in this sector having this property. Then no closed trajectory which encloses the origin can have points in common with  $\Gamma_1$ .

**Proof:** At points of  $\Gamma_1$  where  $F(\rho, \theta) \neq 0$ , the direction field of integral curves is tangent to the radius vector drawn from the origin. Since  $\Phi(\rho, \theta)$  changes sign as  $\Gamma_1$  is crossed, the coordinate  $\theta$  on the integral curves has an extreme value when the integral curve meets  $\Gamma_1$ . Also, since there is no other branch of  $\Phi(\rho, \theta) = 0$  in the sector, on crossing which  $\Phi(\rho, \theta)$  changes sign, a trajectory entering this sector will always remain within it. This proves the lemma.

**Corollary:** If the graph of the equation  $\Phi(\rho, \theta) = 0$  has a unique branch in the sector  $\theta_1 \leq \theta \leq \theta_2$ , with one end at the origin and the other extending to infinity, and if  $\Phi(\rho, \theta)$  changes sign on crossing this branch, then the origin is acyclic.

Let us now consider the system of differential equations

$$\begin{aligned} dx/dt &= a_{10}x + a_{01}y + a_{20}x^2 + a_{11}xy + a_{02}y^2 \\ dy/dt &= b_{10}x + b_{01}y + b_{20}x^2 + b_{11}xy + b_{02}y^2 \end{aligned} \quad (3)$$

with constant coefficients. The following theorem holds.

**Theorem:** In the system (3) a singular point of nodal type cannot be enclosed by a closed trajectory.

**Proof:** Let us assume that the origin is a singular point of nodal type, and let us introduce a system of polar coordinates. The transformed system has the form

$$\begin{aligned} d\rho/dt &= \rho[a_{10}\cos^2\theta + (a_{01} + b_{10})\cos\theta\sin\theta + \\ &\quad b_{01}\sin^2\theta + \rho[a_{20}\cos^3\theta + (a_{11} + b_{20})\cos^2\theta\sin\theta + \\ &\quad (a_{02} + b_{11})\cos\theta\sin^2\theta + b_{02}\sin^3\theta]] \\ d\theta/dt &= b_{10}\cos^2\theta + (b_{01} - a_{10})\cos\theta\sin\theta - \\ &\quad a_{01}\sin^2\theta + \rho[b_{20}\cos^3\theta + (b_{11} - a_{20})\cos^2\theta\sin\theta + \\ &\quad (b_{02} - a_{11})\cos\theta\sin^2\theta - \\ &\quad a_{02}\sin^3\theta] \equiv f_2(\theta) + \rho f_3(\theta) \equiv \Phi(\rho, \theta) \end{aligned} \quad (4)$$

Since  $x = 0, y = 0$  is a singular point of nodal type,  $f_2(\theta)$  either vanishes identically or has a real zero,  $\theta_1$ . The equation  $f_3(\theta) = 0$ , since it is a homogeneous polynomial of third degree in  $\sin\theta$  and  $\cos\theta$ , always has at least one real root. If  $f_2(\theta) \equiv 0$ , if  $f_3(\theta) \equiv 0$ , or if  $f_2(\theta_1) = 0$  and  $f_3(\theta_1) = 0$ , then the conclusion of the theorem follows from the fact that the system (3) has at least one integral curve which is a straight line through the origin. It thus remains to consider the case when the equations  $f_2(\theta) = 0$  and  $f_3(\theta) = 0$  have no common root. Clearly we can always find a sector  $(\theta_1, \theta_2)$  such that  $f_2(\theta_1) = 0$  and  $f_3(\theta_2) = 0$ , and such that neither  $f_2(\theta)$  nor  $f_3(\theta)$  vanishes for any other value of  $\theta$  in this sector. Consider the curve  $\rho = -f_2(\theta)/f_3(\theta)$  in this sector. Without loss of generality we may assume that the function  $-f_2(\theta)/f_3(\theta)$  is positive within the sector, since if it were negative, we would consider the opposite sector  $(\theta_1 \pm \pi, \theta_2 \pm \pi)$ , in which it would be positive. Clearly  $\rho = -f_2(\theta)/f_3(\theta)$  represents a branch of the graph of  $\Phi(\rho, \theta) = 0$  for the system (4), which has one end at the origin and the other extending to infinity, wholly contained in a sector less than  $2\pi$ . As for the variation of sign of  $\Phi(\rho, \theta)$  on crossing this branch, this follows from the linearity of  $\Phi(\rho, \theta)$  in  $\rho$ . There are no other branches of  $\Phi(\rho, \theta) = 0$  in this sector. Thus the conditions of the corollary to the lemma are satisfied. The theorem is proved.

Taking into account the fact that the system (3) can have no more than two singular points of "focus" type,<sup>1</sup> we obtain, as a corollary to the theorem just proved, a known result:<sup>2</sup> If the system (3) has three limit cycles,  $L_c^i (i = 1, 2, 3)$ , then it is impossible for  $I_c^i \cap I_c^j = \emptyset$  for  $i \neq j, j = 1, 2, 3$ . ( $I_c^k$  denotes the domain enclosed by  $L_c^k$ .)

We mention in conclusion that this problem was proposed in the seminar of N. P. Erugin and Yu. S. Bogdenov. The forementioned example was presented there.

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## References

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# Convergence of a Generalized Interpolation Polynomial

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**I**N his paper<sup>1</sup> S. N. Bernshtein proved that the generalized polynomial which deviates least from zero on the interval  $[-1, 1]$  among all generalized polynomials of the form

$$\frac{x^n + b_1x^{n-1} + \dots + b_n}{\prod_{i=1}^n \left(1 - \frac{x}{a_i}\right)} = A_0 + \sum_{k=1}^n \frac{A_k}{a_k - x}$$

has the representation

$$r_n^*(x) = C \cos(n\theta + \psi) \quad (1)$$

where  $C$  is a definite constant

$$\cos\theta = x \quad \psi = 2 \sum_{k=1}^n \alpha_k$$

$$\cos\alpha_k = \frac{\rho_k - x}{\sqrt{2\rho_k(a_k - x)}}$$

$$\sin\alpha_k = \frac{\sqrt{1 - x^2}}{\sqrt{2\rho_k(a_k - x)}}$$

$$a_k = \frac{1}{2} \left( \rho_k + \frac{1}{\rho_k} \right)$$

$$|\rho_k| > 1 \quad k = 1, 2, \dots, n$$

We consider generalized interpolation polynomials on the interval  $[-1, 1]$  with respect to the system of Chebyshev functions

$$\frac{1}{a_1 - x}, \quad \frac{1}{a_2 - x}, \quad \dots, \quad \frac{1}{a_n - x} \quad (2)$$

$$|a_i| > 1 \quad i = 1, 2, \dots, n$$

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